Chapter 8

Billiards

The dynamics that we have the best intuitive grasp on, and find easiest to grapple with both numerically and conceptually, is the dynamics of billiards. For billiards, discrete time is altogether natural; a particle moving through a billiard suffers a sequence of instantaneous kicks, and executes simple motion inbetween, so there is no need to contrive a Poincaré section. We have already used this system in sect. 1.3 as the intuitively most accessible example of chaos. Here we define billiard dynamics more precisely, anticipating the applications to come.

8.1 Billiard dynamics

A billiard is defined by a connected region $Q \subset \mathbb{R}^D$, with boundary $\partial Q \subset \mathbb{R}^{D-1}$ separating Q from its complement $\mathbb{R}^D \setminus Q$. The region Q can consist of one compact, finite volume component (in which case the billiard phase space is bounded, as for the stadium billiard figure 8.1), or can be infinite in extent, with its complement $\mathbb{R}^D \setminus Q$ consisting of one or several finite or infinite volume components (in which case the phase space is open, as for the 3-disk pinball game figure 1.3).

In what follows we shall most often restrict our attention to *planar billiards*. A point particle of mass m and momentum $p_n = mv_n$ moves freely within the billiard, along a straight line, until it encounters the boundary. There it reflects specularly (*specular* = mirrorlike), with no change in the tangential component of momentum, and instantaneous reversal of the momentum component normal to the boundary,

$$p' = p - 2(p \cdot \hat{n})\hat{n} \,, \qquad (8.1)$$

with $\hat{n}$ the unit vector normal to the boundary ∂Q at the collision point. The angle of incidence equals the angle of reflection, as illustrated in figure 8.2. A billiard is a Hamiltonian system with a $2D$-dimensional phase space $x = (q, p)$ and potential $V(q) = 0$ for $q \in Q$, $V(q) = \infty$ for $q \in \partial Q$.

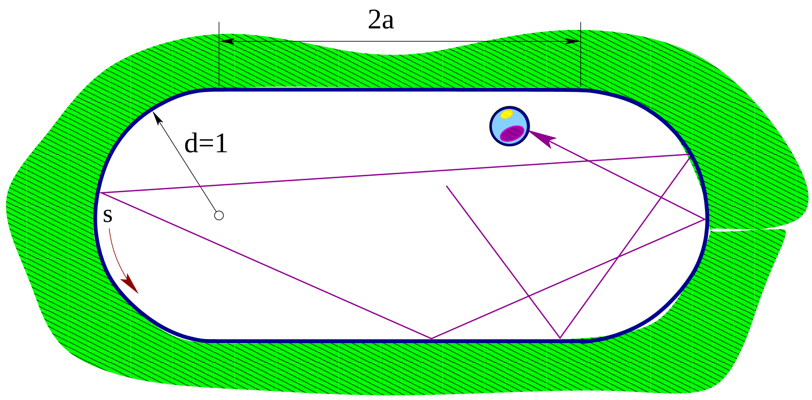

Figure 8.1: The stadium billiard is a 2-dimensional domain bounded by two semi-circles of radius $d = 1$ connected by two straight walls of length $2a$. At the points where the straight walls meet the semi-circles, the curvature of the border changes discontinuously; these are the only singular points of the flow. The length a is the only parameter.

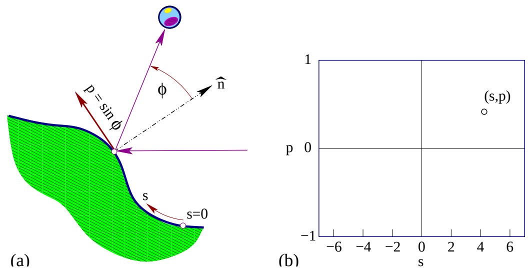

Figure 8.2: (a) A planar billiard trajectory is fixed by specifying the perimeter length parametrized by s and the outgoing trajectory angle ϕ, both measured counterclockwise with respect to the outward normal $\hat{n}$. (b) The Birkhoff phase space coordinate pair (s, p) fully specifies the trajectory, where $p = |p| \sin\phi$ is the momentum component tangential to the boundary (and we set $|p| = 1$ whenever convenient).

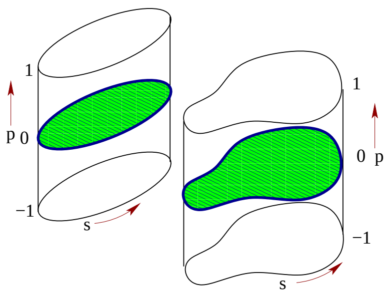

Figure 8.3: In $D = 2$ the billiard Poincaré section is a cylinder, with the parallel momentum p ranging over $p \in \{-1, 1\}$, and with the s coordinate is cyclic along each connected component of ∂Q. The rectangle figure 8.2 (b) is such cylinder unfolded, with periodic boundary conditions glueing together the left and the right edge of the rectangle.

A billiard flow has a natural Poincaré section defined by Birkhoff coordinates s_n, the arc length position of the nth bounce measured along the billiard boundary, and $p_n = |p| \sin \phi_n$, the momentum component parallel to the boundary, where ϕ_n is the angle between the outgoing trajectory and the normal to the boundary. We measure both the arc length s, and the parallel momentum p counterclockwise relative to the outward normal (see figure 8.2 as well as figure 1.8 (a)). In $D = 2$, the Poincaré section is a cylinder (topologically an annulus), figure 8.3, where the parallel momentum p ranges for $-|p|$ to $|p|$, and the s coordinate is cyclic along each connected component of ∂Q. The volume in the full phase space is preserved by the Liouville theorem (6.32). The Birkhoff coordinates $x = (s, p) \in \mathcal{P}$, are the natural choice, because with them the the Poincaré return map preserves the phase space volume in the (s, p) parameterize Poincaré section (a perfectly good coordinate set (s, ϕ) does not do that).

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Without loss of generality we set $m = |v| = |p| = 1$. Poincaré section condition eliminates one dimension, and the energy conservation $|p| = 1$ eliminates another, so the Poincaré section return map P is $(2D - 2)$-dimensional.

The dynamics is given by the Poincaré return map

$$P : (s_n, p_n) \mapsto (s_{n+1}, p_{n+1}) \tag{8.2}$$

from the nth collision to the $(n + 1)$st collision. The discrete time dynamics map P is equivalent to the Hamiltonian flow (6.1) in the sense that both describe the same full trajectory. Let t_n denate the instant of nth collision. Then the position of the pinball $\in Q$ at time $t_n + \tau \leq t_{n+1}$ is given by $2D - 2$ Poincaré section coordinates $(s_n, p_n) \in \mathcal{P}$ together with τ, the distance reached by the pinball along the nth section of its trajectory.

Example 8.1 3-disk game of pinball: *In case of bounces off a circular disk, the position coordinate $s = r\theta$ is given by angle $\theta \in [0, 2\pi]$. For example, for the 3-disk game of pinball of figure 1.7 and figure 1.8 we have two types of collisions:*

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$$P_0 : \begin{cases} \phi' = -\phi + 2 \arcsin p \\ p' = -p + \frac{a}{R} \sin \phi' \end{cases} \qquad \textit{back-reflection} \tag{8.3}$$

$$P_1 : \begin{cases} \phi' = \phi - 2 \arcsin p + 2\pi/3 \\ p' = p - \frac{a}{R} \sin \phi' \end{cases} \qquad \textit{reflect to 3rd disk}. \tag{8.4}$$

Here a = radius of a disk, and R = center-to-center separation. Actually, as in this example we are computing intersections of circles and straight lines, nothing more than high-school geometry is required. There is no need to compute $\arcsin$'s either - one only needs to compute a square root per each reflection, and the simulations can be very fast.

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Trajectory of the pinball in the 3-disk billiard is generated by a series of P_0's and P_1's. At each step on has to check whether the trajectory intersects the desired disk (and no disk inbetween). With minor modifications, the above formulas are valid for any smooth billiard as long as we replace a by the local curvature of the boundary at the point of collision.

8.2 Stability of billiards

We turn next to the question of local stability of discrete time billiard systems. Infinitesimal equations of variations (4.2) do not apply, but the multiplicative structure (4.44) of the finite-time fundamental matrices does. As they are more physical than most maps studied by dynamicists, let us work out the billiard stability in some detail.

On the face of it, a plane billiard phase space is 4-dimensional. However, one dimension can be eliminated by energy conservation, and the other by the fact that the magnitude of the velocity is constant. We shall now show how going to a local frame of motion leads to a [2×2] fundamental matrix.

Consider a 2-dimensional billiard with phase space coordinates $x = (q_1, q_2, p_1, p_2)$. Let t_k be the instant of the kth collision of the pinball with the billiard boundary, and $t_k^{\pm} = t_k \pm \epsilon$, ϵ positive and infinitesimal. With the mass and the velocity equal to 1, the momentum direction can be specified by angle θ: $x = (q_1, q_2, \sin\theta, \cos\theta)$. Now parametrize the 2-$d$ neighborhood of a trajectory segment by $\delta x = (\delta z, \delta\theta)$, where

$$\delta z = \delta q_1 \cos\theta - \delta q_2 \sin\theta \,, \tag{8.5}$$

$\delta\theta$ is the variation in the direction of the pinball motion. Due to energy conservation, there is no need to keep track of $\delta q_\parallel$, variation along the flow, as that remains constant. $(\delta q_1, \delta q_2)$ is the coordinate variation transverse to the kth segment of the flow. From the Hamilton's equations of motion for a free particle, $dq_i/dt = p_i$, $dp_i/dt = 0$, we obtain the equations of motion (4.1) for the linearized neighborhood

$$\frac{d}{dt}\delta\theta = 0, \quad \frac{d}{dt}\delta z = \delta\theta \,. \tag{8.6}$$

Let $\delta\theta_k = \delta\theta(t_k^+)$ and $\delta z_k = \delta z(t_k^+)$ be the local coordinates immediately after the kth collision, and $\delta\theta_k^- = \delta\theta(t_k^-)$, $\delta z_k^- = \delta z(t_k^-)$ immediately before.

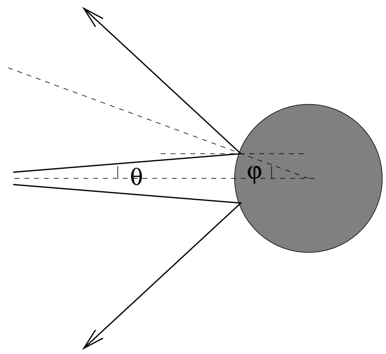

Figure 8.4: Defocusing of a beam of nearby trajectories at a billiard collision. (A. Wirzba)

Integrating the free flight from t^+_{k-1} to t^-_k we obtain

$$\begin{aligned} \delta z^-_k &= \delta z_{k-1} + \tau_k \delta\theta_{k-1}\,, \qquad \tau_k = t_k - t_{k-1} \\ \delta\theta^-_k &= \delta\theta_{k-1}\,, \end{aligned} \tag{8.7}$$

and the stability matrix (4.43) for the kth free flight segment is

$$M_T(x_k) = \begin{pmatrix} 1 & \tau_k \\ 0 & 1 \end{pmatrix} . \tag{8.8}$$

At incidence angle ϕ_k (the angle between the outgoing particle and the outgoing normal to the billiard edge), the incoming transverse variation δz^-_k projects onto an arc on the billiard boundary of length $\delta z^-_k / \cos\phi_k$. The corresponding incidence angle variation $\delta\phi_k = \delta z^-_k/\rho_k \cos\phi_k$, $\rho_k =$ local radius of curvature, increases the angular spread to

$$\begin{aligned} \delta z_k &= -\delta z^-_k \\ \delta\theta_k &= -\,\delta\theta^-_k - \frac{2}{\rho_k \cos\phi_k} \delta z^-_k\,, \end{aligned} \tag{8.9}$$

so the fundamental matrix associated with the reflection is

$$M_R(x_k) = -\begin{pmatrix} 1 & 0 \\ r_k & 1 \end{pmatrix} , \qquad r_k = \frac{2}{\rho_k \cos\phi_k} . \tag{8.10}$$

The full fundamental matrix for n_p consecutive bounces describes a beam of trajectories defocused by M_T along the free flight (the τ_k terms below) and defocused/refocused at reflections by M_R (the r_k terms below)

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$$M_p = (-1)^{n_p} \prod_{k=n_p}^{1} \begin{pmatrix} 1 & \tau_k \\ 0 & 1 \end{pmatrix} \begin{pmatrix} 1 & 0 \\ r_k & 1 \end{pmatrix} , \tag{8.11}$$

where τ_k is the flight time of the kth free-flight segment of the orbit, $r_k = 2/\rho_k \cos\phi_k$ is the defocusing due to the kth reflection, and ρ_k is the radius of curvature of the billiard boundary at the kth scattering point (for our 3-disk game of pinball, $\rho = 1$). As the billiard dynamics is phase space volume preserving, $\det M = 1$, and the eigenvalues are given by (6.22).

This is still another example of the fundamental matrix chain rule (4.51) for discrete time systems, rather similar to the Hénon map stability (4.52). Stability of every flight segment or reflection taken alone is a shear with two unit eigenvalues,

$$\det M_T = \det \begin{pmatrix} 1 & \tau_k \\ 0 & 1 \end{pmatrix}, \quad \det M_R = \det \begin{pmatrix} 1 & 0 \\ r_k & 1 \end{pmatrix}, \tag{8.12}$$

but acting in concert in the intervowen sequence (8.11) they can lead to a hyperbolic deformation of the infinitesimal neighborhood of a billiard trajectory.

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As a concrete application, consider the 3-disk pinball system of sect. 1.3. Analytic expressions for the lengths and eigenvalues of $\overline{0}$, $\overline{1}$ and $\overline{10}$ cycles follow from elementary geometrical considerations. Longer cycles require numerical evaluation by methods such as those described in chapter 12.

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chapter 12

Résumé

A particulary natural application of the Poincaré section method is the reduction of a billiard flow to a boundary-to-boundary return map.

Commentary

Remark 8.1 Billiards. The 3-disk game of pinball is to chaotic dynamics what a pendulum is to integrable systems; the simplest physical example that captures the essence of chaos. Another contender for the title of the 'harmonic oscillator of chaos' is the baker's map which is used as the red thread through Ott's introduction to chaotic dynamics [1.12]. The baker's map is the simplest reversible dynamical system which is hyperbolic and has positive entropy. We will not have much use for the baker's map here, as due to its piecewise linearity it is so nongeneric that it misses all of the subtleties of cycle expansions curvature corrections that will be central to this treatise.

☞ chapter 18

That the 3-disk game of pinball is a quintessential example of deterministic chaos appears to have been first noted by B. Eckhardt [8.1]. The model was studied in depth classically, semiclassically and quantum mechanically by P. Gaspard and S.A. Rice [8.2], and used by P. Cvitanović and B. Eckhardt [8.3] to demonstrate applicability of cycle expansions to quantum mechanical problems. It has been used to study the higher order $\hbar$ corrections to the Gutzwiller quantization by P. Gaspard and D. Alonso Ramirez [8.4], construct semiclassical evolution operators and entire spectral determinants by P. Cvitanović and G. Vattay [8.5], and incorporate the diffraction effects into the periodic orbit theory by G. Vattay, A. Wirzba and P.E. Rosenqvist [8.6]. The full quantum mechanics and semiclassics of scattering systems is developed here in the 3-disk scattering context in chapter 33. Gaspard's monograph [1.8], which we warmly recommend, utilizes the 3-disk system in much more depth than will be attained here. For further links check `ChaosBook.org`.

A pinball game does miss a number of important aspects of chaotic dynamics: generic bifurcations in smooth flows, the interplay between regions of stability and regions of chaos, intermittency phenomena, and the renormalization theory of the 'border of order' between these regions. To study these we shall have to face up to much harder challenge, dynamics of smooth flows.

Nevertheless, pinball scattering is relevant to smooth potentials. The game of pinball may be thought of as the infinite potential wall limit of a smooth potential, and pinball symbolic dynamics can serve as a *covering* symbolic dynamics in smooth potentials. One may start with the infinite wall limit and adiabatically relax an unstable cycle onto the corresponding one for the potential under investigation. If things go well, the cycle will remain unstable and isolated, no new orbits (unaccounted for by the pinball symbolic dynamics) will be born, and the lost orbits will be accounted for by a set of pruning rules. The validity of this adiabatic approach has to be checked carefully in each application, as things can easily go wrong; for example, near a bifurcation the same naive symbol string assignments can refer to a whole island of distinct periodic orbits.

☞ sect. 26.1

Remark 8.2 Stability analysis. The chapter 1 of Gaspard monograph [1.8] is recommended reading if you are interested in Hamiltonian flows, and billiards in particular. A. Wirzba has generalized the stability analysis of sect. 8.2 to scattering off 3-dimensional spheres (follow the links in `ChaosBook.org/extras`). A clear discussion of linear stability for the general d-dimensional case is given in Gaspard [1.8], sect. 1.4.

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